# **Problem Solutions**

# **Problem 1**

## Part A

For the light traveling to Reflector  $R_1$ , the light that reaches the reflector must be going "upwind" relative to the interferometer, as shown. Therefore, its speed relative to the interferometer is

$$\sqrt{c^2 - v^2}$$

The light travels a total distance 2L, so the time is

$$t_1 = \frac{2L}{\sqrt{c^2 - v^2}}$$
(1.1)

For the light traveling to Reflector  $R_2$ , when it is traveling towards the reflector its speed is c + v relative to the interferometer, and when it is returning to the beam splitter its speed is c - v. Thus the total time is:

$$t_2 = \frac{L}{c+v} + \frac{L}{c-v} = \frac{2Lc}{c^2 - v^2}$$
(1.2)

From Equations 1.1 and 1.2 it is fairly simple to show the Equation A.1 is true.

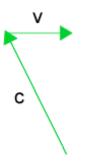
### Part B

The difference in times is

$$t_2 - t_1 = \frac{2Lc}{c^2 - v^2} - \frac{2L}{\sqrt{c^2 - v^2}} = 2L \left[ \frac{c}{c^2 - v^2} - \frac{1}{\sqrt{c^2 - v^2}} \right]$$
(1.3)

Evaluating this can be a bit of a challenge for cheap calculators because the two terms in the square brackets have almost the same value. Here is a form which may work better for your calculator

$$t_2 - t_1 = \frac{2L}{\sqrt{c^2 - v^2}} \left[ \frac{c}{\sqrt{c^2 - v^2}} - 1 \right] = 3.33 \times 10^{-16} \,\mathrm{s} \tag{1.4}$$



Alternatively we can use the fact that  $(1-x)^n \approx 1-nx$  for small x to show that the time difference is approximately

$$t_2 - t_1 \approx \frac{Lv^2}{c^3} \approx 3.33 \times 10^{-16} \,\mathrm{s}$$

In terms of the phase difference, this is

$$\Delta \phi = 2\pi \frac{t_2 - t_1}{T} = 2\pi f (t_2 - t_1) = 2\pi \times 0.200 \text{ rads}$$
(1.5)

## Part C

This is just all of the above except that the labels *1* and *2* are interchanged. Thus we can immediately write the answer:

$$\Delta \phi = -2\pi \times 0.200 \, \text{rads}$$

# Part D

Recall that when the phase difference is  $\pi$  we have complete destructive interference. Here we have gone from a phase difference of 0.400  $\pi$  through 0 (constructive interference) to -0.400  $\pi$  radians. This should be observable.

## **Problem 2**

If the period of the source is *T* relative to Pablo, then the wavelength of the wave relative to him is

$$\lambda = (c+v)T \tag{2.1}$$

But this period is the time dilated value. If the period of the source in a frame stationary relative to the source is  $T_0$ , then

$$T = \gamma T_0 = \frac{1}{\sqrt{1 - v^2 / c^2}} T_0 \tag{2.2}$$

The wavelength of the wave in a frame stationary relative to the source,  $\lambda_0$ , is related to this period by

$$T_0 = \frac{\lambda_0}{c} \tag{2.3}$$

Thus Equation 2.1 becomes

$$\lambda = (c+v)\frac{1}{\sqrt{1-v^2/c^2}}\frac{\lambda_0}{c} = \sqrt{\frac{1+v/c}{1-v/c}}\lambda_0$$
(2.4)

# **Problem 3**

Let's start with the value of zero. This is just the value if a signal propagating at c travels from Event 1 to Event 2. Although this may be obvious, let's do the math

$$c^2 \Delta t^2 - \Delta x^2 = 0 \tag{3.1}$$

Divide by  $\Delta t^2$ .

$$c^2 - \left(\frac{\Delta x}{\Delta t}\right)^2 = 0 \tag{3.2}$$

But  $\frac{\Delta x}{\Delta t}$  is just the speed of a signal propagating from Event 1 to Event 2, v, so:

$$v = c \tag{3.3}$$

If the interval squared is greater than 0, Eqn 3.2 becomes:

$$c^2 - \left(\frac{\Delta x}{\Delta t}\right)^2 > 0 \tag{3.4}$$

So

$$v < c \tag{3.5}$$

This makes sense. If the two events occur at the same place at different times, the interval squared is positive.

If the interval squared is less than 0, we get

$$v > c \tag{3.6}$$

This also makes sense. If the two events occur simultaneously at different places, the interval squared is less than 0, and a signal that connects them would have to move at infinite speed.

### **Problem 4**

#### Part 1

If  $u_{\text{Lou}} = v$ , then *f* is 1. This is reasonable: the particle doesn't move relative to the train and the light meets it at the rear of the train.

#### Part 2

If  $u_{\text{Lou}} = c$ , then *f* is 0. This too is reasonable: the "race" ends in a tie.

#### Part 3

For Sue, the speed of light is (c - v) when it is traveling to the front of the car and (c + v) after it is reflected. Thus the equivalent of Eqn C.4 is

$$u_{Sue}(t_{1,Sue} + t_{2,Sue}) = (c - v)t_{1,Sue} - (c + v)t_{2,Sue}$$
(4.1)

Therefore

$$\frac{t_{2,Sue}}{t_{1,Sue}} = \frac{(c-v) - u_{Sue}}{(c+v) + u_{Sue}}$$
(4.2)

For Sue, the equivalents of Eqns. C.5 and C.6 are

$$(c-v)t_{1,Sue} = L_{Sue}$$

$$(c+v)t_{2,Sue} = fL_{Sue}$$
(4.3)

Eliminate  $L_{Sue}$ , solve for f and use Eqn. 4.2 and we get

$$f = \frac{(c+v)(c-v-u_{Sue})}{(c-v)(c+v+u_{Sue})}$$
(4.4)

Equating this to Eqn. C.7 we end up with

$$u_{Sue} = u_{Lou} - v \tag{4.5}$$

This is the "common sense" result gotten by Galileo.

# Problem 5

Energy is the time component of the vector 4-momentum. This answer is just from looking at the structure of the equation.