# Lecture 24

# **Gaussian Beam Optics**

### **Perturbation Analysis**

- Many of the examples we have treated have phase which
  - varies fast in the DOP (z)
  - varies slowly as a function of the other co-ordinates (x,y)
- Fresnel diffraction varies quadratically in the x,y direction
  - $(1/r)\exp(ikr) \rightarrow (1/z)\exp(ikz)\exp(ik(x^2+y^2)/(2z))$
  - separates the fast-changing z-direction from the slow changing x,y direction
- We can generalise!
- write the wave as u(x,y,z)exp(ikz) subst in wave equation

$$\nabla_T^2 u + \frac{\delta^2 u}{\delta z^2} + 2ik\frac{\delta u}{\delta z} = 0$$

•  $\nabla_{T}$  is the 2-D grad function

### **Paraxial Wave Equation**

- If we now assume that the function
  - varies slowly in z on the scale of a wavelength |∂u/∂z| « k|u|
  - that it is smooth higher order differentials can be ignored

$$\nabla_T^2 u + \frac{\delta^2 u}{\delta z^2} + 2ik\frac{\delta u}{\delta z} \rightarrow \nabla_T^2 u + 2ik\frac{\delta u}{\delta z} = 0$$

- The slowly varying envelope approximation (SVEA) leads to the paraxial wave equation
- Note that u(x,y,z)=(1/z)exp(ik(x<sup>2</sup>+y<sup>2</sup>)/(2z)) is a solution of the above equation
- Trouble with that solution is that it has infinite extent is there a similar solution which has finite extent - ie looks something like a pencil beam?

### A Gaussian Beam

# Yes, there is.... $(1/z)\exp(ik(x^2+y^2)/(2z)) \rightarrow$ $u_{00}'(x,y,z) = (1/(z-iz_0))\exp(ik(x^2+y^2)/(2(z-iz_0)))$ $z \rightarrow z - iz_0$

Normalise over x, y

$$\begin{aligned} u_{00}(x,y,z) &= \sqrt{\frac{kz_0}{\pi}} \frac{1}{z - iz_0} \exp\left(\frac{ik[x^2 + y^2]}{2(z - iz_0)}\right) \\ &= \sqrt{\frac{2}{\pi}} \frac{1}{w} e^{-i\varphi} \exp\left(-\frac{[x^2 + y^2]}{w^2}\right) \exp\left(\frac{ik[x^2 + y^2]}{2R}\right) \end{aligned}$$

• 
$$w^2(z) = w_0^2(1 + z^2/z_0^2),$$
  $w_0^2 = 2z_0/k$   
•  $R = (z^2 + z_0^2)/z,$   $\tan \phi = z/z_0$ 

### **Properties of the Gaussian Beam**

$$U_{00}(x,y,z) = \sqrt{\frac{2}{\pi}} \frac{1}{w} e^{-i\varphi} \exp\left(-\frac{r^2}{w^2}\right) \exp\left(\frac{ikr^2}{2R}\right)$$

- Fundamental Gaussian Beam Solution
- One parameter  $z_0$  describes the beam
- Circularly symmetric function of r
- Gaussian extent transversely w is the e<sup>-1</sup> point of amplitude
- at z = 0,  $w = w_0$  and is the minimum extent of beam
- at  $z = z_0$ ,  $w = \sqrt{2} w_0 z_0$  is called the *confocal parameter*

### **Properties of the Gaussian Beam**

$$u_{00}(x,y,z) = \sqrt{\frac{2}{\pi} \frac{1}{w}} e^{-i\varphi} \exp\left(-\frac{r^2}{w^2}\right) \exp\left(\frac{ikr^2}{2R}\right)$$

- What is the significance of R?
- R is the radius of curvature of the surfaces of constant phase

■ as 
$$z \rightarrow 0$$
,  $R \rightarrow \infty$ 

- as  $z \to \infty$ ,  $R \to z$
- What is the significance of  $\varphi$ ?
  - φ is related to the velocity of surfaces of constant phase
  - Remember this isn't a plane wave and doesn't even satisfy the full wave equation



### **Any Other Solutions?**

Yes, it can be shown that...

$$U_{l,m}(X,Y,Z) =$$

$$\frac{C_{l,m}}{W} H_{l}\left[\frac{x\sqrt{2}}{W}\right] H_{m}\left[\frac{y\sqrt{2}}{W}\right] e^{-i(l+m+1)\phi} \exp\left(-\frac{r^{2}}{W^{2}}\right) \exp\left(\frac{ikr^{2}}{2R}\right)$$

- Where H<sub>m</sub>(x) is a Hermite polynomial
- C<sub>I,m</sub> is a normalisation constant
- I and m are integers
- AKA TEM<sub>Im</sub>
- Phase speed is

$$k_{eff} = k - (l + m + 1) \frac{Z_0}{Z^2 + Z_0^2}$$

### **Paraxial Optics With Gaussian Beams**

- A Gaussian beam is a function of  $z_0$
- Gaussian beam at any position is a function of  $q = z i z_0$
- A Gaussian beam is locally a spherical wave with radius of curvature R
- Consider a Gaussian beam on a thin lens
- The size of the beam doesn't change on either side of the lens
- The curvature of the beam is altered by the lens
- IT CAN BE SHOWN THAT
  - The transformation of the parameter q is given by
  - $q^{(Aq + B)}/(Cq + D)$
  - where A,B,C,D are the elements of the 2x2 matrix for paraxial ray optical systems
  - Notice that this is NOT a matrix equation

### **Transformations of Gaussian Beams**

- Stated that u(x,y,z)=(1/q)exp(ik(x<sup>2</sup>+y<sup>2</sup>)/(2q)) is a solution to the paraxial wave equation
- Where q = z i z<sub>0</sub>, z is the distance in the DOP and z<sub>0</sub> describes the Gaussian beam (Gaussian beam parameter)
- By inspection of wave equation solution show q(z) propagates as:
- $1/q(z) = 1/R(z) + i \lambda/(\pi n w(z)^2)$
- Far field propagate like spherical wave R(z)->z, z-> $\infty$

### **Free Space Gaussian Beam**

- Consider a Gaussian beam propagating in free space
- Reference frame for Gaussian beams defines waist at z=0
- At plane  $z=z_1$  we can write  $q_1 = z_1 i z_0$
- At plane  $z=z_2$  we can write  $q_2 = z_2 i z_0 = q_1 + (z_2 z_1)$
- Plane shift corresponds to change in real component of q
- Can calculate beam width and curvature at new position as



# Gaussian Beams and Optics

- Consider a thin lens with an incident Gaussian Beam
- Spot size does not change as beam passes through lens
- But, lens changes wave front (earlier Fraunhoffer lecture)
   T(x,y) = exp(-ik(x<sup>2</sup>+y<sup>2</sup>) / 2f)
- Apply this directly to wave equation solution at lens boundary => effect is to change R->R'

$$\frac{ik(x^2+y^2)}{2R'} = \frac{ik(x^2+y^2)}{2R} - \frac{ik(x^2+y^2)}{2f}$$

### **Gaussian Beams and Optics II**

- We can simplify this as 1/R' = 1/R 1/f
- So, 1/q' = 1/q 1/f [the spot size does not change]
- We can therefore write the lens transformation as

$$q' = q / (-q/f + 1)$$

- Equivalent to defining a new origin and Gaussian beam parameter for the propagating wave
- As previously outlined, the general transformation is

$$q^{} = (Aq + B)/(Cq + D)$$

- Known as a fractional linear or Möbius transformation
- To confuse matters further the transformation is also called the ABCD law and is written as a pseudo-matrix [A B; C D]
- THIS HAS **NOTHING** TO DO WITH MATRIX ALGEBRA!

### **Common Transformations**

- Free Space Transformation becomes [1 z<sub>2</sub>-z<sub>1</sub>; 0 1]
- Thin lens with focal length f becomes [1 0; -1/f 1]
- Note that converging lenses have +ve f by convention
- Series of optical components can be represented by result of successive transformations in q space

$$M_{series} = M_1 \dots M_{last}$$

 Multiple transformations are applied in an order determined by the beam propagation

### **Gaussian Beam Summary**

- A Gaussian beam is the paraxial limit solution to a wave equation
- Equivalent to a point source shifted by an imaginary amount iz<sub>0</sub>
- With z<sub>0</sub> = 0 paraxial solution is just portion of spherical wave emanating from the origin
- q = z i z<sub>0</sub> measures complex distance from reference plane (waist of Guassian beam)
- Transforms in q-space can be used to model effect of optics on a Gaussian beam

### **Applications of Gaussian Beam Optics**

- Results are remarkably similar to ray tracing but equations include wave behaviour of light
- Gaussian beam does not behave quite like a plane wave and consequently equivalent calculations give slightly different answers
- Laser people routinely carry out these calculations
- Simple example next lecture...

### **Gaussian Beam Optics: Example**



**Question:** A He-Ne laser operating in Gaussian mode has a divergence of 1mR and with a beam waist at the output of 0.4mm. What is the diffraction limited spot size that we can achieve with a lens f=+2cm located 1m from the beam waist?